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# Energy Conditions and Statefinder Diagnostic in $f(R, T)$ Gravity with an Anisotropic Background

**Original Research  
Article**

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## Abstract

In our study, we explored the properties of a spatially homogeneous and anisotropic Bianchi type  $VI_0$  Universe. Our investigation centered on integrating cosmic domain walls into the  $f(R, T)$  theory of gravitation, initially proposed by Harko et al. (2011). To tackle the field equations, we employed the relationship between the expansion scalar ( $\theta$ ) and the shear scalar ( $\sigma$ ). Our analysis encompassed both the dynamic and cosmological aspects of the Universe. By comparing our findings to the  $\Lambda$ CDM model, specifically focusing on the evolution of the jerk parameter, we found a striking agreement between the two models. A noteworthy discovery was the verification of accelerated expansion in our described model, consistent with the prevailing observational data. Finally, we examine the energy condition criteria and determine that the violation of the Strong Energy Condition (SEC), while the Null Energy Condition (NEC), Weak Energy Condition (WEC) and Dominant Energy Condition (DEC) continue to meet the requirements for positivity.

*Keywords:* Bianchi type  $VI_0$  cosmological model,  $f(R, T)$  gravity, Domain wall, Power law  
Astronomy and Astrophysics: 53C25; 83C05; 57N16

## 1 Introduction

In contemporary cosmology, a significant revelation is that the Universe is undergoing both expansion and acceleration. This late-time accelerated expansion has been confirmed through research on high red-shift supernovae (Riess et al., 1998, Perlmutter et al., 1999, Bennet et al., 2003). The investigation of the Universe, which seems pervaded by dark energy, has captivated numerous scientists. Research findings have made it evident that the Universe is predominantly influenced

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by a unique energy form with negative pressure commonly termed dark energy. The cosmological constant, as explored in the Padmanabhan's work (Padmanabhan, 2003), plays a pivotal role in deciphering the nature of this dark energy. Evidence from diverse sources like the Cosmic Microwave Background Radiation (CMBR) and supernova surveys has unveiled that the Universe's energy composition comprises about 4% regular baryonic matter, 22% dark matter, and 74% dark energy (Riess et al., 2004, Eisenstein et al., 2005, Astier et al., 2006, Spergel et al., 2007). Recent times have witnessed the proposition of several alterations to the theory of general relativity (GR) to provide a natural gravitational framework for understanding dark energy. In the quest to explain the Universe's late-time acceleration, researchers are actively exploring alternative avenues. Among these, the  $f(R)$  theory of gravity stands out as a suitable candidate due to its cosmological implications. Introduced by replacing the Einstein-Hilbert action of GR with a generalized function of the Ricci scalar  $R$  (Nojiri and Odintsov, 2007, Multamaki and Vilja, 2006, 2007, Shamir, 2010),  $f(R)$  gravity embodies the amalgamation of early-time inflation and the Universe's late-time acceleration.

Studies on the Bianchi type- $V I_0$  cosmological model indicate its potential convergence towards isotropy (Adhav et al., 2011). Analyzing this model further, it becomes evident that its accelerated expansion is attributed to a negative barotropic equation of state (Yadav et al., 2012). With the passage of time, the deceleration of the Bianchi type- $V I_0$  Universe increases gradually, eventually reaching a constant value (Shaikh and Katore, 2016). Scholars (Bali and Kumari, 2017, Satish and Venkateswarlu, 2019) have explored various aspects of this model, including its shearing, non-rotational, and expanding nature. In the context of Lyra geometry, researchers (Pradhan et al., 2007) have investigated solutions for a bulk viscous plane symmetric Universe and highlighted the absence of big bang singularities. Another examination of solutions involving a plane symmetric cosmological model with a domain wall demonstrates the presence of radiation (Pawar et al., 2009). A significant discovery involves solutions for a non-static Bianchi type-III cosmological model with domain walls, both in the presence and absence of a magnetic field, within the framework of general relativity (Adhav et al., 2009). Lastly, in the presence of string and domain walls associated with quarks, researchers (Sahoo and Mishra, 2013) have identified a vacuum kink model that notably lacks a singularity at  $r = 2k$ .

Modifications to General Relativity (GR) involved coupling matter and geometry through a Lagrangian dependent on the stress-energy tensor trace ( $T$ ) and the Ricci scalar ( $R$ ). The  $f(R, T)$  gravity field equations are derived from the Hilbert-Einstein principle (Harko et al., 2011). Despite setting up energy density and pressure for dark components, equilibrium thermodynamics isn't achievable in  $f(R, T)$  gravitation. In thermal equilibrium, photons and non-photons follow the generalized second law of thermodynamics (Sharif and Zubair, 2012). Analyzing the Kantowski-Sachs bulk viscous Universe showed pressure, density, viscosity, Hubble parameter tend toward zero for high cosmic time (Khade et al., 2018). Investigating different  $f(R, T)$  gravity models (with  $n = 0$  and  $n \neq 0$ ) revealed an expanding, shearing, non-rotating, accelerating Universe (Hasmani and Al-Haysah, 2019). Under dark energy's influence, expansion occurs; negative deceleration and positive Hubble parameter imply exponential expansion and acceleration (Pawar et al., 2019). The imaginary dark energy form acts as quintessence field when  $\omega = \frac{1}{3}$ , explaining accelerated expansion (Islam et al., 2019). With a stiff fluid, the Bianchi type- $V$  Universe shows isotropy in  $f(R, T)$  gravity (Patil et al., 2020). Using domain walls as fractal cosmology's matter source, the flat Friedmann-Robertson-Walker model contracts and accelerates (Pawar et al., 2020). Linear and nonlinear  $f(R, T)$  gravity leads to accelerated expansion like  $\Lambda$ CDM model (Sahoo et al., 2021). Bianchi type models depict shearing, non-rotation, accelerated expansion, approaching isotropy as cosmic time increases (Chaubey and Shukla, 2017). Some scholars utilized the first-order formalism in Horndeski cosmology to simplify the equations of motion for Horndeski gravity, as well as  $(\square)$  and  $(\square, \square)$  theories. Our numerical solutions aligned well with current observations, particularly the Hubble parameter's behavior over redshift (Santos et al. 2019). Some researchers have explored a brane cosmology scenario involving an inflating 3D domain wall within a five-dimensional Minkowski space, alongside a stack of  $N$  parallel domain walls. They analyzed aspects of inflation theory governed by a scalar field that drives the

inflation (Neves et al., 2020). Some scholars have explored BPS domain walls within Horndeski gravity, specifically focusing on AdS vacua and discovered that, unlike double sine-Gordon models, the system can connect UV and IR fixed points without bypassing intermediate vacua (Santos and Brito, 2024).

In the context of  $f(R, T)$  gravity, where a perfect fluid serves as the energy source, it has been observed that the geometry of a Bianchi Type  $VI_0$  Universe remains undisturbed. However, there is a slight alteration in the matter distribution, as discussed in reference (Rao and Nilima, 2013). For Bianchi Type III and Kantowski-Sachs Universes, predictions indicate a future collapse attributed to domain walls within the  $f(R, T)$  theory. This collapse contributes to the model's stability, ensuring the absence of singularities in the Universe (Katore and Hatkar, 2016). The impact of bulk viscosity is noted on pressure and the equation of state parameter, yet, it does not affect the density of the domain wall itself (Mahanta et al., 2018). The presence of domain walls leads to an expanding nature of Riemannian space-time, particularly for extended periods. This expansion aligns with observations related to Type Ia supernovae (Shaikh and Wankhade, 2018). In the initial stages, the Universe exhibits expansion with a finite volume that further increases with time. As time ( $t$ ) progresses, the model tends towards isotropy at  $t = 0$  (Pawar et al., 2021). By delving into the study of a Bianchi type-V Universe within the framework of  $f(R, T)$  gravity, and considering the presence of both domain walls and quark matter, it is established that pressure and density experience growth as redshift ( $z$ ) increases. This scenario points towards the existence of a Big Bang singularity and confirms the Universe's accelerated expansion, resembling observations related to Type Ia Supernovae (Maurya et al., 2020). Moreover, researchers have explored the Plane Symmetry cosmological model within the context of  $f(R, T)$  gravity with interacting fields. The outcomes of this study align closely with recent observational data, revealing an expanding Universe (Pawar and Mapari, 2022). The collaborative research endeavors inspire a deeper exploration of the spatially homogeneous anisotropic Bianchi Type  $VI_0$  Universe. In this context, we consider a universe that contains a cosmic domain wall, which acts as the source of energy within the framework of the  $f(R, T)$  theory of gravity.

## 2 Equations of $f(R, T)$ Gravity

The action principle for  $f(R, T)$  modified theory is given by

$$S = \frac{1}{16\pi G} \int f(R, T) \sqrt{-g} d^4x + \int L_m \sqrt{-g} d^4x \quad (2.1)$$

where,  $f(R, T)$  is a function of Ricci scalar  $R$  and  $T$  be trace of the stress of energy tensor and  $L_m$  is Lagrangian of the matter.

The stress energy of the matter is

$$T_{ij} = -\frac{2}{\sqrt{-g}} \frac{\partial(\sqrt{-g})}{\partial g^{ij}} L_m, \quad \Theta_{ij} = -2T_{ij} - p g_{ij} \quad (2.2)$$

By varying the action principle with respect to  $g_{ij}$ , the corresponding field equations of  $f(R, T)$  gravity are obtained as,

$$f_R(R, T) R_{ij} - \frac{1}{2} f(R, T) g_{ij} + \left( g_{ij} \nabla^i \nabla_j - \nabla_i \nabla_j \right) f(R, T) = 8\pi T_{ij} - f_T(R, T) T_{ij} - f_T(R, T) \Theta_{ij} \quad (2.3)$$

Where,  $f_R = \frac{\delta f(R, T)}{\delta R}$ ,  $f_T = \frac{\delta f(R, T)}{\delta T}$  and  $\Theta_{ij} = g^{\alpha\beta} \frac{\delta T_{\alpha\beta}}{\delta g^{\alpha\beta}}$

Here, covariant derivative is represented by  $\nabla$  and the energy momentum tensor  $T_{ij}$  emerges from the Lagrangian  $L_m$ . By assuming the function  $f(R, T) = f(R)$ , the equation (2.3) is reduced to field equations of  $f(R)$  gravity.

In this paper, we assume  $f(R, T)$  is of the form,

$$f(R, T) = f_1(R) + f_2(T) \quad (2.4)$$

where

$$f_1(R) = \lambda R \quad \text{and} \quad f_2(T) = \lambda T \quad (2.5)$$

in which,  $\lambda$  is arbitrary constant.

Using equation (2.2) and (2.4), the equation (2.3) becomes

$$f_1(R) R_{ij} - \frac{1}{2} f(R, T) g_{ij} + (g_{ij} \nabla^i \nabla_i - \nabla_i \nabla_j) f_1(R) = 8\pi T_{ij} - f_2(T) T_{ij} - f_2(T) [-2T_{ij} - p g_{ij}]$$

$$G_{ij} = R_{ij} - \frac{1}{2} R g_{ij} = \kappa T_{ij} + \Lambda g_{ij} \quad (2.6)$$

### 3 Metric and Field Equations

The spatially homogeneous and anisotropic Bianchi type  $VI_0$  line element can be written in the form as,

$$ds^2 = dt^2 - A^2 dx^2 - B^2 e^{2x} dy^2 - C^2 e^{-2x} dz^2 \quad (3.1)$$

Where,  $A$ ,  $B$  and  $C$  are the functions of time  $t$  only.

The energy momentum tensor for domain wall is given by

$$T_{ij} = (g_{ij} + \omega_i \omega_j) \rho - p \omega_i \omega_j \quad (3.2)$$

Where,  $p$  and  $\rho$  are the pressure and density of the fluid respectively and  $\omega_i = (0, 0, 0, 1)$  is four velocity vector satisfying  $\omega_i \omega^j = 0$  and  $\omega_i \omega^i = -1$ .

With the help of co-moving coordinates system and equation (3.2), the Einstein field equations for the cosmological model (3.1) are given by

$$\frac{\ddot{B}}{B} + \frac{\ddot{C}}{C} + \frac{\dot{B}\dot{C}}{BC} + \frac{1}{A^2} = \kappa\rho - \Lambda \quad (3.3)$$

$$\frac{\ddot{A}}{A} + \frac{\ddot{C}}{C} + \frac{\dot{A}\dot{C}}{AC} - \frac{1}{A^2} = \kappa\rho - \Lambda \quad (3.4)$$

$$\frac{\ddot{A}}{A} + \frac{\ddot{B}}{B} + \frac{\dot{A}\dot{B}}{AB} - \frac{1}{A^2} = \kappa\rho - \Lambda \quad (3.5)$$

$$\frac{\dot{A}\dot{B}}{AB} + \frac{\dot{B}\dot{C}}{BC} + \frac{\dot{A}\dot{C}}{AC} - \frac{1}{A^2} = -\kappa p - \Lambda \quad (3.6)$$

$$\frac{\dot{B}}{B} - \frac{\dot{C}}{C} = 0 \quad (3.7)$$

The overhead dot ( $\dot{\phantom{x}}$ ) denotes the derivative with respect to time  $t$ .

Integrating equation (3.7) we have,

$$B = lC$$

Where  $l$  is a constant of integration.

Without loss of generality, here we choose  $l = 1$ , we have

$$B = C \quad (3.8)$$

Using equation (3.8), equations (3.3)-(3.6) reduced to

$$2 \frac{\ddot{B}}{B} + \frac{\dot{B}^2}{B^2} + \frac{1}{A^2} = \kappa\rho - \Lambda \quad (3.9)$$

$$\frac{\ddot{A}}{A} + \frac{\ddot{B}}{B} + \frac{\dot{A}\dot{B}}{AB} - \frac{1}{A^2} = \kappa\rho - \Lambda \quad (3.10)$$

$$2 \frac{\dot{A}\dot{B}}{AB} + \frac{\dot{B}^2}{B^2} - \frac{1}{A^2} = -\kappa p - \Lambda \quad (3.11)$$

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## 4 Solution of Field Equations

There are three linearly independent equations with five unknowns  $A$ ,  $B$ ,  $p$ ,  $\rho$  and  $\Lambda$ . Therefore to solve this system of equations we assume that the expansion scalar is proportional to shear scalar. This condition leads to

$$A = B^n, \quad n \neq 0 \quad (4.1)$$

From equations (3.9) and (3.10) we have

$$\frac{\ddot{B}}{B} + \frac{\dot{B}^2}{B^2} + \frac{2}{A^2} - \frac{\ddot{A}}{A} - \frac{\dot{A}\dot{B}}{AB} = 0 \quad (4.2)$$

Using equation (4.1) in (4.2) we get

$$B = [n(k_1 t + k_2)]^{\frac{1}{n}} \quad (4.3)$$

Where,  $k_1$  &  $k_2$  are the constants of integration and

$$k_1^2 = \frac{1}{n-1}, \quad n \neq 1$$

From equation (3.8), (4.1) and (4.3) we have

$$A = n(k_1 t + k_2) \quad \& \quad C = n(k_1 t + k_2)^{\frac{1}{n}} \quad (4.4)$$

Putting the values of  $A$ ,  $B$  and  $C$ , in equation (3.1) yields

$$ds^2 = dt^2 - n^2(k_1 t + k_2)^2 dx^2 - [n(k_1 t + k_2)]^{\frac{2}{n}} [e^{2x} dy^2 - e^{-2x} dz^2] \quad (4.5)$$

The Volume ( $V$ ) is

$$V = ABC$$

$$V = [n(k_1 t + k_2)]^{\frac{n+2}{n}} \quad (4.6)$$

The Hubble Parameter ( $H$ ) is

$$H = \frac{1}{3}(H_1 + H_2 + H_3)$$

Where,  $H_1 = \frac{\dot{A}}{A}$ ,  $H_2 = \frac{\dot{B}}{B}$  and  $H_3 = \frac{\dot{C}}{C}$

$$H = \frac{1}{3} \frac{k_1(n+2)}{n(k_1 t + k_2)} \quad (4.7)$$

The Scalar expansion ( $\theta$ ) is

$$\theta = 3H$$

$$\theta = \frac{k_1(n+2)}{n(k_1 t + k_2)} \quad (4.8)$$

The deceleration parameter ( $q$ ) is

$$q = \frac{d}{dt} \left( \frac{1}{H} \right) - 1$$

$$q = \frac{2(n-1)}{(n+2)}, \quad (4.9)$$

The Anisotropic parameter ( $A_m$ ) is

$$A_m = \frac{1}{3} \sum_{i=1}^3 \left( \frac{H_i - H}{H} \right)^2$$

$$A_m = 2 \left( \frac{n-1}{n+2} \right)^2 \quad (4.10)$$

The Shear Scalar ( $\sigma^2$ ) is

$$\sigma^2 = \frac{3}{2} A_m H^2$$

$$\sigma^2 = \frac{1}{6} \left( \frac{k_1(n-1)}{n(k_1 t + k_2)} \right)^2 \quad (4.11)$$

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## 5 Dynamical Properties of model

The equation of state (EoS) for cosmic domain wall is

$$p = -\frac{2}{3}\rho \quad (5.1)$$

From equation (3.10) and (3.11) we have

$$\frac{\ddot{A}}{A} + \frac{\ddot{B}}{B} - \frac{\dot{A}\dot{B}}{AB} - \frac{\dot{B}^2}{B^2} = \kappa(p + \rho) \quad (5.2)$$

Using equation (5.1), equation (5.2) yields The density ( $\rho$ ) is

$$\rho = \frac{2}{\kappa n} \left( \frac{k_1}{(k_1 t + k_2)} \right)^2 \quad (5.3)$$

Form equation (5.1) we have The Pressure ( $p$ ) is

$$p = -\frac{4}{3\kappa n} \left( \frac{k_1}{(k_1 t + k_2)} \right)^2 \quad (5.4)$$

The cosmological constant ( $\Lambda$ ) is

$$\Lambda = \frac{7}{3\kappa n} \left( \frac{k_1}{(k_1 t + k_2)} \right)^2 \quad (5.5)$$

## Jerk Paramater

In cosmology, the term "jerk parameter" refers to a dimensionless quantity that characterizes the rate of change of the acceleration of the expansion of the Universe. This parameter plays a crucial role in understanding the dynamics of the cosmos and helps us investigate the nature of dark energy, a mysterious force driving the accelerated expansion of the Universe.

The expansion of the Universe was initially thought to be slowing down due to the gravitational pull of matter, which is the dominant component in the Universe. However, observations of distant supernovae in the late 1990s provided compelling evidence that the expansion is actually accelerating. This unexpected discovery led to the proposal of dark energy, a hypothetical form of energy that counteracts the attractive force of gravity, as the driving force behind this acceleration.

The jerk parameter, denoted as " $j$ ", is a higher-order cosmological parameter that comes into play when studying the evolution of the cosmic scale factor,  $a(t)$ , which describes how the Universe expands over time. It is related to the third derivative of the scale factor with respect to cosmic time,  $a(t)$ . is defined and discussed in the ref. (Blandford et al., 2004, Rapetti et al., 2007, Chiba and Nakamura, 1998, Visser, 2004, Luongo, 2013)

The cosmic scale factor is given by

$$a(t) = V^{1/3} \quad (5.6)$$

Using equation (4.6) in equation (5.6) yields

$$a(t) = [n(k_1 t + k_2)]^{\frac{n+2}{3n}} \quad (5.7)$$

$$j(t) = \frac{1}{aH^3} \frac{d^3}{dt^3} (a) \quad (5.8)$$

From equation (5.7) and equation (5.8), we have

$$j(t) = \frac{2(1-n)(2-5n)}{(n+2)^2} \quad (5.9)$$

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## 6 Energy Conditions

Energy conditions encompass a collection of principles and inequalities within the framework of general relativity. They serve as essential tools for characterizing the behavior of energy-momentum tensors in the fabric of spacetime. These conditions assume a pivotal role in our comprehension of the curvature of spacetime and its implications for extraordinary phenomena, including concepts such as wormholes or warp drives. In the following discussion, we will delve into an overview of several fundamental energy conditions.

1. **Strong Energy Condition (SEC):** The Strong Energy Condition posits that gravitational forces should always be attractive. In terms of the energy-momentum tensor, it is expressed as:  
 $\rho + 3P \geq 0$
2. **Null Energy Condition (NEC):** The Null Energy Condition states that for any null vector  $\mu^\alpha$ , the following inequality must hold:  $\rho + P \geq 0$
3. **Weak Energy Condition (WEC):** The Weak Energy Condition (WEC) asserts that the energy density observed by any observer must always remain non-negative, implying that it cannot be negative under any circumstances i.e.  $\rho + P \geq 0$  and  $\rho \geq 0$
4. **Dominant Energy Condition (DEC):** The Dominant Energy Condition (DEC) serves as a criterion that ensures the energy density observed by any observer must remain non-negative, implying that  
 $\rho \geq |P|$  i.e.  $\rho - P \geq 0$  and  $\rho + P \geq 0$ .

from equations (5.3) and (5.4) we found that

**SEC :**

$$\rho + 3P = -\frac{2}{\kappa n} \left( \frac{k_1}{(k_1 t + k_2)} \right)^2 \leq 0 \quad (6.1)$$

**NEC:**

$$\rho + P = \frac{2}{3\kappa n} \left( \frac{k_1}{(k_1 t + k_2)} \right)^2 \geq 0 \quad (6.2)$$

and

$$\rho = \frac{2}{\kappa n} \left( \frac{k_1}{(k_1 t + k_2)} \right)^2 \geq 0 \quad (6.3)$$

**WEC**

$$\rho + P = \frac{2}{3\kappa n} \left( \frac{k_1}{(k_1 t + k_2)} \right)^2 \geq 0 \quad (6.4)$$

**DEC**

$$\rho - P = \frac{10}{3\kappa n} \left( \frac{k_1}{(k_1 t + k_2)} \right)^2 \geq 0 \quad (6.5)$$

## 7 Statefinder Parameters

The statefinder parameters is a set of dimensionless parameters introduced by Sahni & Starobinsky (2000) to provide a more insightful understanding of the cosmic dynamics and the nature of dark energy. These parameters, denoted as  $\{r, s\}$ , are used to diagnose the evolution and characteristics of the Universe and its components. They are particularly useful for distinguishing between different cosmological models, such as the Cold Dark Matter with a Cosmological Constant ( $\Lambda$ CDM) model and the Standard Cold Dark Matter (SCDM) model, by identifying unique fixed points  $\{r, s\} = \{1, 0\}$  and  $\{r, s\} = \{1, 1\}$  respectively correspond to each model's properties. Statefinder parameters offer

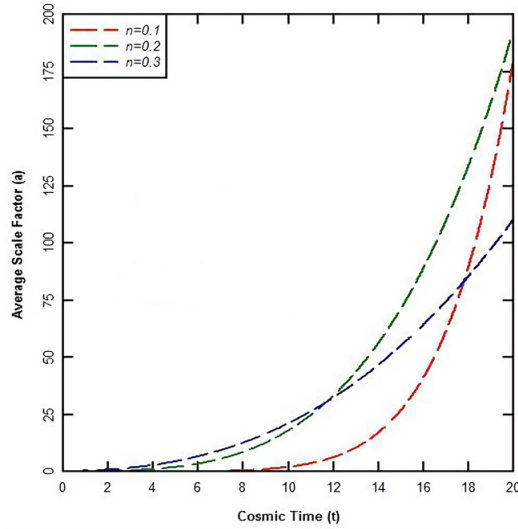


Figure 1: Variation of Average scale factor ( $a$ ) against cosmic time ( $t$ ) for  $n = 0.1, n = 0.2, n = 0.3$

a geometric perspective on the Universe's expansion and are a valuable tool in cosmological studies. The state finder parameter  $r$  and  $s$  are defined as follows:

$$r = \frac{1}{aH^3} \frac{d^3}{dt^3} (a), \quad s = \frac{r-1}{3(q-\frac{1}{2})} \quad (7.1)$$

Using equations (4.9) and (5.7), above equations yields

$$\{r, s\} = \left\{ \frac{2(1-n)(2-5n)}{(n+2)^2}, \frac{19n^2-32n+4}{9(n+2)(n-2)} \right\} \quad (7.2)$$

## 8 Observations

From the figures we have observed that,

- The increasing graph of the average scale factor and volume of the Universe as depicted in Figure 1 and Figure 2 serve as compelling indicators of the Universe's expansion, a pivotal concept in our comprehension of the cosmos. These findings align with empirical observations like the redshift of light emanating from distant galaxies, providing strong support for the Big Bang theory and the notion of a Universe in a state of continuous expansion.
- It is observed that the Hubble parameter  $H$  is decreasing function of cosmic time ( $t$ ) in the positive region (Figure 3). The positive value of  $H$  confirmed about the expansion of the Universe. Initially, the rate of expansion is faster but later on it slows down as time increases and it will be zero for large time ( $t$ ) (Figure 4).
- We found that shear scalar is a diminishing function of cosmic time ( $t$ ) (Figure 5). The present model is not shear free except for  $n = 1$ .

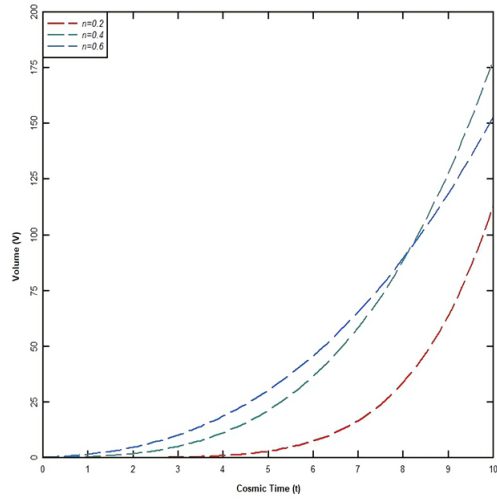


Figure 2: Variation of Volume (V) against cosmic time (t) for  $n = 0.2, n = 0.4, n = 0.6$

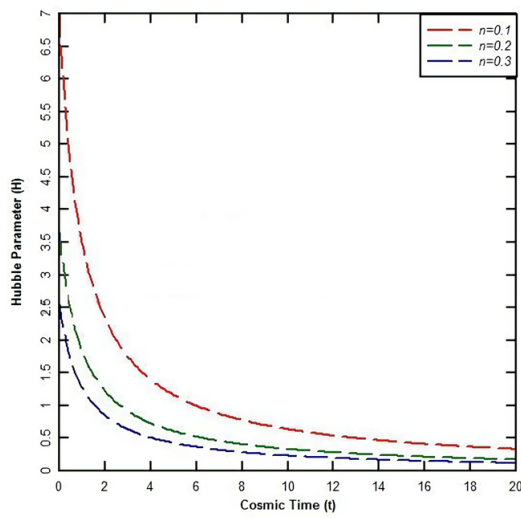


Figure 3: Variation of Hubble parameter (H) against cosmic time (t) for  $n = 0.1, n = 0.2, n = 0.3$

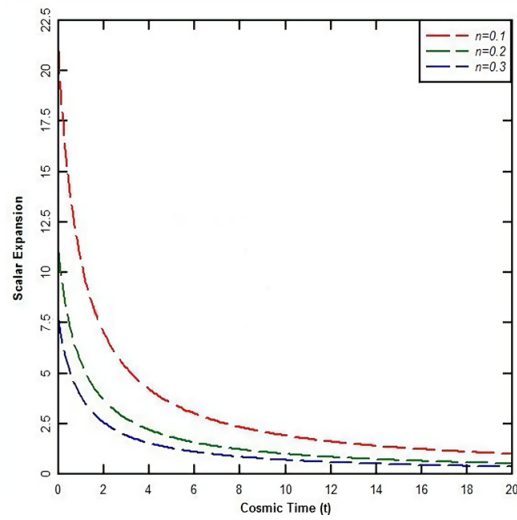


Figure 4: Variation of Scalar expansion ( $\theta$ ) against cosmic time ( $t$ ) for  $n = 0.1, n = 0.2, n = 0.3$

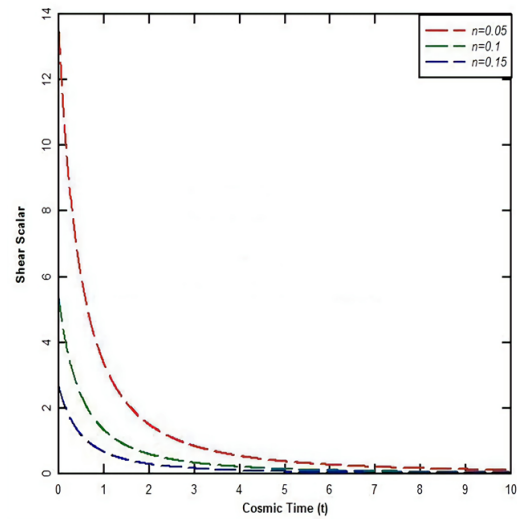


Figure 5: Variation of Shear scalar ( $\sigma$ ) against cosmic time ( $t$ ) for  $n = 0.05, n = 0.1, n = 0.15$

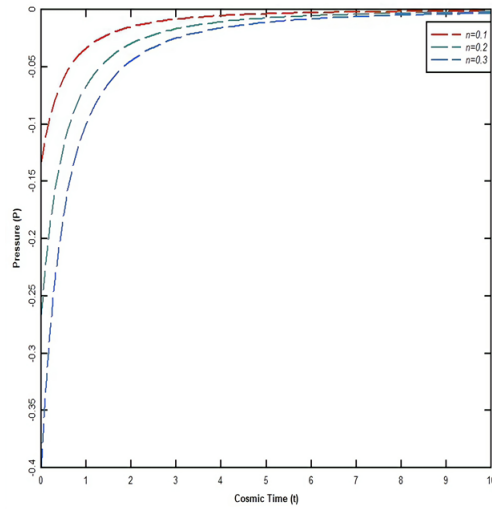


Figure 6: Variation of Pressure ( $P$ ) against cosmic time ( $t$ ) for  $n = 0.1, n = 0.2, n = 0.3$

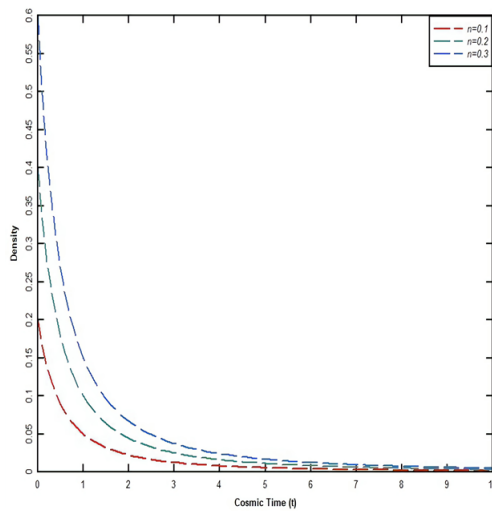


Figure 7: Variation of Density ( $\rho$ ) against cosmic time ( $t$ ) for  $n = 0.1, n = 0.2, n = 0.3$

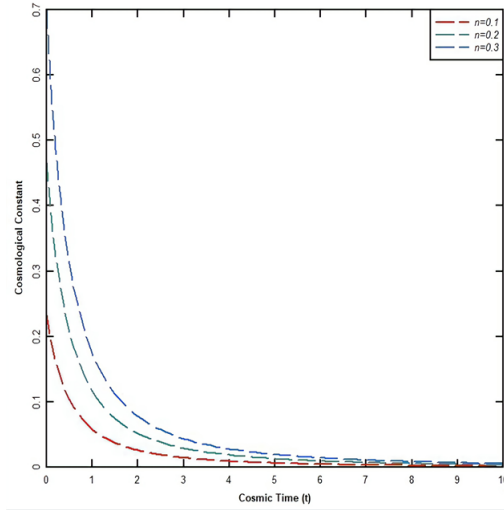


Figure 8: Variation of Cosmological constant ( $\Lambda$ ) against cosmic time (t) for  $n = 0.1, n = 0.2, n = 0.3$

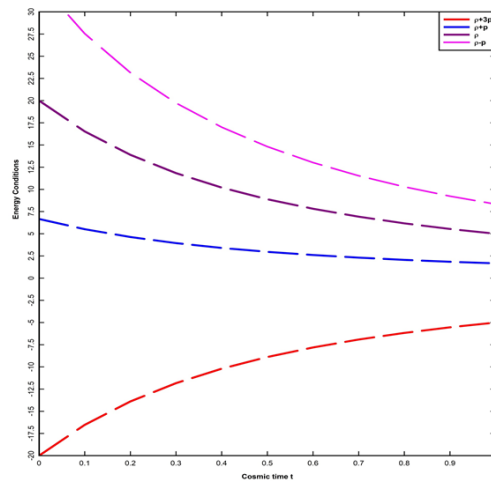


Figure 9: Variation of Energy Conditions (EC) against cosmic time (t) for  $n = 0.1$

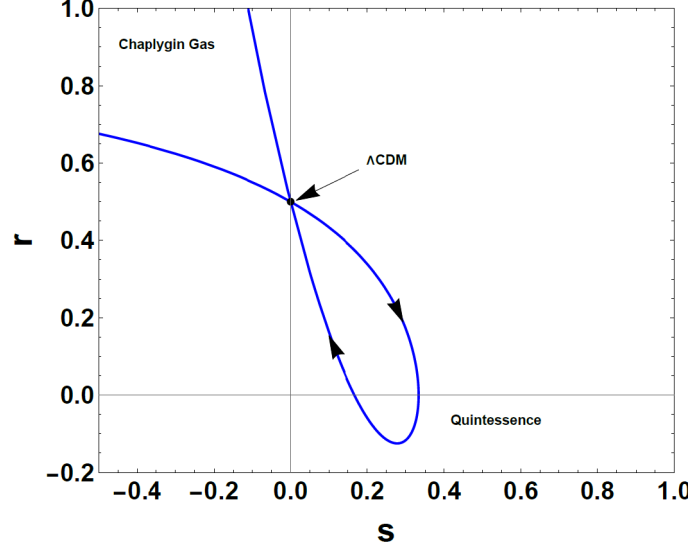


Figure 10: Evolution of  $r - s$  trajectory

- The pressure vary from large negative value to small negative value (Figure 6) and it tends to zero for large time  $t$  ( $t \rightarrow \infty$ ). A negative nature of the pressure shows that the existence of dark energy. Density is the decreasing function of cosmic time ( $t$ ) (Figure 7). It approaches to zero for infinite time.
- We found the variable cosmological constant which is decreasing function of cosmic time ( $t$ ) (Figure 8).
- Through observations, it has been established that the Weak Energy Condition (WEC), Null Energy Condition (NEC), and Dominant Energy Condition (DEC) are all satisfied during the expansion of the Universe. However, it is noteworthy that the Strong Energy Condition (SEC) is found to be violated, as illustrated in Figure 9.
- Figure 10 illustrates the progression trajectory of statefinder parameters corresponding to the specified  $f(R, T)$  gravity model.

Here all above quantities i.e. average scale factor ( $a$ ), volume ( $V$ ), Hubble parameter ( $H$ ), scalar expansion ( $\theta$ ), shear scalar ( $\sigma$ ), pressure ( $P$ ), density ( $\rho$ ) and cosmological constant ( $\Lambda$ ) are in arbitrary units.

## 9 Concluding Remarks

In the framework of the  $f(R, T)$  theory of gravity, we derived exact solutions to the field equations of an anisotropic Bianchi type- $V I_0$  cosmological model of the Universe with cosmic domain wall as an energy source. The metric potentials are finite at initial epoch, increasing as time increases and model does not have initial singularity. We have investigated the profile of cosmological and dynamical parameters in the framework of  $f(R, T)$  gravity. In the present described model, we have used theoretical model as a cosmological constant for describing the nature of the dark energy and noticed that the cosmological constant is positive and dependent function of cosmic time ( $t$ ). It indicates

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towards the observations of Supernovae Ia experiment (Riess, et al., 1998). We have found the existence of dark energy because of negative pressure. Also, we observed that the Universe is expanding as the Hubble parameter ( $H$ ) is positive and the volume ( $V$ ) of the Universe is increasing with cosmic time ( $t$ ). Also, the rate of expansion of the Universe is decreasing with the cosmic time ( $t$ ). It is in good agreement with the observations of Pawar et al., 2021. The deceleration parameter (DP) has a negative value for  $0 < n < 1$  which points to the accelerating phase of the Universe. From equation (4.10) and (4.11), it is clear that our Universe is verified as anisotropic and not shear free throughout the evolution of the Universe for ( $n \neq 1$ ). Equation of state (EoS) plays a major role to explore the nature of the Universe and we have used EoS for cosmic domain wall ( $p = -\frac{2}{3}\rho$ ) to describe the dynamical parameters of the model. We have observed the pressure and density are finite at initial time and decreasing functions of cosmic time ( $t$ ). Also, we have explored the jerk parameter and compared with  $\Lambda$ CDM model where the value of jerk is given by  $j(t) = 1$ . In the present discussed model we found  $j(t) = 1$  for  $n = 2$  from equation (5.9), which represents the exact  $\Lambda$ CDM model. Also, as  $n \rightarrow 2^-$  ( $n < 2$ ), the model approaches to  $\Lambda$ CDM ( $j(t) < 1$ ) and as  $n \rightarrow 2^+$  ( $n > 2$ ), then it shows departures from  $\Lambda$ CDM model. We found the cosmological constant  $\Lambda$  is a function of cosmic time  $t$  and agreed with the observations of numerous researchers (Poplawski, 2006, Zel'dovich, 1968, Linde, 1974). We have observed the variable cosmological constant  $\Lambda$  gives a positive value, which indicates the expansion of the Universe is accelerating which reassembles with observations of the present Universe's behavior. The statefinder parameter trajectory traverses through the  $\Lambda$ CDM fixed point and subsequently extends into the quintessence region ( $r < 1, s > 0$ ) before ultimately reaching the Chaplygin gas region ( $r > 1, s < 0$ ) as time unfolds into the distant future. Finally, the Raychaudhuri equations demonstrate the satisfaction of the Null Energy Condition (NEC), Weak Energy Condition (WEC), and Dominant Energy Condition (DEC) signifies non-negative energy density, affirming the model's physical viability and the ongoing accelerated expansion of the universe. However, the Strong Energy Condition (SEC) is found to be violated. This SEC violation aligns with previous findings by (Sahoo et al., 2020, Patil et al. 2023), indicating an accelerated expansion of the Universe.

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## References

- [1] Adhav, K. S., Bansod, A. S., Munde, S. L. and Nakwal, R. G., 2011, *Astrophys. Space Sci.*, 332, 497-502 <http://dx.doi.org/10.1007/s10509-010-0519-3>
- [2] Adhav, K. S., Raut, V. B., Thakare, R. S. and Kale, C. B., 2009, *Int. J.Theo. Phys.*, 48(8), 2290-2296. <http://dx.doi.org/10.1007/s10773-010-0530-z>
- [3] Astier, P., et al., 2006, *Astron. Astrophys.*, 447, 31-48. <http://dx.doi.org/10.1051/0004-6361:20054185>
- [4] Bali, R. and Kumari, P., 2017, *Advances in Astrophysics*, 2, 2. <http://dx.doi.org/10.22606/adap.2017.22001>

- 
- [5] Bennet, C. L. et al., 2003, *Astrophys. J., Suppl.*, 148 (1). <http://dx.doi.org/10.1086/377253>  
bibitemr6 Blandford, R. D. , Amin, M. , Baltz, E. A. , Mandel, K. and Marshall, P. J., 2004, *ASP Conf. Ser.*, 339(27) <https://doi.org/10.48550/arXiv.astro-ph/0408279>
- [6] Capozziello, S., et al., 2018, *Phys. Lett. B* 781 099
- [7] Chaubey, R. and Shukla, A. K., 2017, *Pramana J. Phys.*, 88 65. <http://dx.doi.org/10.1007/s12043-017-1371-6>
- [8] Chiba, T. and Nakamura T., 1998, *Prog. Theor. Phys.* 100 1077 <https://doi.org/10.1143/PTP.100.1077>
- [9] Eisenstein, D. J., et al., 2005, *Astrophys. J.*, 633, 560-576. <http://dx.doi.org/10.1086/466512>
- [10] Harko, T., Lobo, F. S. N., Nojiri, S. and Odintsov, S. D., 2011, *Phys. Rev., D* 84, 024020. <http://dx.doi.org/10.1103/PhysRevD.84.024020>
- [11] Hasmani, A. H. and Al-Haysah, A. M., 2019, *JMSM*, 2(3), 163-175. <http://dx.doi.org/10.33187/jmsm.435340>
- [12] Islam, S., Kumar, P., Khadekar, G. S. and Das, T. K., 2019, *J. Phys.: Conference Series*, 1258, 012026. <http://dx.doi.org/10.1088/1742-6596/1258/1/012026>
- [13] Katore, S. D. and Hatkar, S. P., 2016, *Prog. Theor. Exp. Phys.* 033E01. <https://doi.org/10.1093/ptep/ptw009>
- [14] Khade, P. P., Wasnik, A. P. and Kandalkar, S. P., 2018, *GJSFR*, 18, 1.
- [15] Linde, A. D., 1974, *JETP Lett.*, 19, 183.
- [16] Luongo, O., 2013, *Mod. Phys. Lett. A* 28, 1350080. <https://doi.org/10.1142/S0217732313500806>
- [17] Mahanta, K. L., Bharati, J. K., Lepse, P. V. and Bishi, B. K., 2018, *JETIR* 5, 10. [www.jetir.org/papers/JETIRDQ06042.pdf](http://www.jetir.org/papers/JETIRDQ06042.pdf)
- [18] Maurya, D. C., Pradhan, A. and Dixit, A., 2020, *Int. J. Geom. Meth. Mod. Phys.*, 17(1) 2050014. <https://doi.org/10.1142/S0219887820500140>
- [19] Moraes, P. H. R. S. and Sahoo, P. K., 2017, *Eur. Phys. J. C* 77 480
- [20] Multamaki, T. and Vilja, I., 2006, *Phys. Rev., D* 74, 064022 <http://dx.doi.org/10.1103/PhysRevD.74.064022>
- [21] Multamaki, T. and Vilja, I., 2007, *Phys. Rev., D* 76, 064021 <http://dx.doi.org/10.1103/PhysRevD.76.064021>
- [22] Nojiri, S. and Odintsov, S. D., 2007, *Int. J. Geom. Meth. Mod. Phys.*, 4, 115-146. <http://dx.doi.org/10.1142/S0219887807001928>
- [23] Padmanabham, T., 2003, *Phys. Rept.*, 380, 235-320
- [24] Patil, V. R. , Pawar, D. D. , Pawde, J. L. and Mapari, R. V., 2020, *Ajanta*, 8(1), 11-18.
- [25] Patil, V. R. , Pawde, J. L. , Mapari, R. V. and Bolke, P. A., 2023, *East Eur. J. Phys.* 3, 62, <https://doi.org/10.26565/2312-4334-2023-3-04@V>.

- 
- [26] Patil, V. R., Bolke, P. A., Waghmare, S. K., and Pawde, J. L., 2023, *East Eur. J. Phys.* 3, 53, <https://doi.org/10.26565/2312-4334-2023-3-03>.
- [27] Pawar, D. D. , Mapari, R. V. and Agrawal, P. K., 2019, *J. Astrophys. Astr.*, 40, 2. <http://dx.doi.org/10.1007/s12036-019-9582-5>
- [28] Pawar, D. D. , Raut, D. K. and Patil, W. D., 2020, *IJMPA*, 35(17) 2050072.
- [29] Pawar, D. D. and Mapari, R. V., 2022, *J. Dyn. Sys. Geom. Theo.*, 20(1) 115-136 <https://doi.org/10.1080/1726037X.2022.2079268>
- [30] Pawar, D. D., Bayaskar, S. N. and Patil, V. R., 2009, *Bulg. J. Phys.*, 36, 68–75.
- [31] Pawar, D. D., Mapari, R. V. and Raut, V. M., 2021, *Bulg. J. Phys.* 48, 225–235.
- [32] Perlmutter, S. et al., 1999, *Astrophys. J.*, 517 565-586. <http://dx.doi.org/0.1086/307221>
- [33] Pop-lawski, N. J., 2006, *Phys. Lett. B* 640, 135. <https://doi.org/10.1016/j.physletb.2006.07.056>
- [34] Pradhan, A., Rai, K. and Yadav, A., 2007, *Braz. J. Phys.*, 37, 3B. <http://dx.doi.org/10.1590/S0103-97332007000700003>
- [35] Rao, V. U. M. and Neelima, D., 2013 *Astrophys Space Sci.*, 345, 427–430. <http://dx.doi.org/10.1007/s10509-013-1406-5>
- [36] Rapetti, D. , Allen, S. W. , Amin, M. A. and Blandford, R. D., 2007, *Mon. Not. R. Astron. Soc.* 375, 1510 <https://doi.org/10.1111/j.1365-2966.2006.11419.x>
- [37] Riess, A. G. et al., 1998, *Astron. J.*, 116 1009-1038. <http://dx.doi.org/10.1086/300499>
- [38] Riess, A. G., et al., 2004, *Astrophys. J.*, 607, 665-687. [http://dx.doi.org/10.1016/S0370-1573\(2803\)2900120-0](http://dx.doi.org/10.1016/S0370-1573(2803)2900120-0)
- [39] Sahni, V. and Starobinsky, A. A., 2000, *Int. J. Mod. Phys. D*, 9(4), 373-443.
- [40] Sahoo, P. K. and Mishra, B., 2013, *J. Theor. Appl. Phys.*, 7, 62. <http://dx.doi.org/10.1186/2251-7235-7-62>
- [41] Sahoo, P. K., Mandal, S. and Arora, S., 2021, *Astron. Nachr.*, 342, 89–95. <http://dx.doi.org/10.1002/asna.202113886>
- [42] Sahoo, P. K., Sahoo P. and Bishi, B. K., 2017, *Int. J. Geom. Methods Mod. Phys.* 14 1750097
- [43] Santos, F. F., Neves, R.M.P. and Brito, F. A. *Ad. High Energy Physics*, 2019, 3486805, (8), <https://doi.org/10.1155/2019/3486805>.
- [44] Santos, F. F. and Brito, F. A., 2024, *Phys. Lett. B*, 850, <https://doi.org/10.1016/j.physletb.2024.138543>.
- [45] Neves, R.M.P., Santos, F. F. and Brito, F. A., 2020, *Phys. Lett. B*, 810, <https://doi.org/10.1016/j.physletb.2020.135813>.
- [46] Satish, J. and Venkateswarlu, R., 2019, *Bulg. J. Phys.*, 46, 67–79.
- [47] Shaikh, A. Y. and Katore, S. D , 2016, *Bulg. J. Phys.*, 43,(3) 184–194.

- 
- [48] Shaikh, A.Y. and Wankhade, K. S., 2018, IJSRST 4, 10, 134-14.
- [49] Shamir, M. F., 2010, *Astrophys. Space Sci.*, 330, 183-189 <http://dx.doi.org/10.1007/s10509-010-0371-5>
- [50] Sharif, M. and Zubair, M., 2012, *JCAP*, 3, 028. <http://dx.doi.org/10.1088/1475-7516/2012/03/028>
- [51] Sharif, M. and Fatima, I., 2014, *Astrophys. Space Sci.* 353 259
- [52] Spergel, D. N., et al., 2007, *Astron. Astrophys.*, 447, 31-48. <http://dx.doi.org/10.1086/513700>
- [53] Visser, M., 2004, *Class. Quant. Grav.* 21, 2603. <https://doi.org/10.1088/0264-9381/21/11/006>
- [54] Yadav, P., Chandel, S., Singh, M. K. and Shri Ram, 2012, *GJSFR*, 12, 12, 1
- [55] Zel'dovich, Ya. B., 1968, *Sov. Phys. JETP Lett.* 6, 316. <https://doi.org/10.1070/PU1968v011n03ABEH003927>.