

# Resolution of the standard telegraph equation by the Laplace-Adomian method

## Abstract

In this paper, we research the solution of the standard telegraph equation by the Laplace-Adomian method. The Laplace-Adomian method is based on the combination of Laplace transform and the Adomian decomposition method.

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**Key word:** Telegraph equation, Laplace transform, ADM method.

## 1 Introduction

In this article, we study the general solution of the standard telegraph equations by the method of Laplace-Adomian. The standard telegraph equation is an important equation arises in the propagation of electrical signals along a telegraph line, taking into consideration the inductance, capacitance and conductance of the cable. However the method of Laplace-Adomian is a numerical method based on the combination of the Laplace Transform and Adomian decomposition method.

## 2 The numerical Laplace-Adomian method

The standard telegraph equation is a partial differential equation given by :

$$\partial^2 u \partial x^2 = \alpha \partial^2 u \partial t^2 + \beta \partial u \partial t + \gamma u$$

where  $u = u(t, x)$  is the resistance, and  $\alpha, \beta$  and  $\gamma$  are constants related to the inductance, capacitance and conductance of the cable respectively.

Let us consider the following functional equation:

$$\begin{cases} \partial^2 u \partial x^2 & = \alpha \partial^2 u \partial t^2 + \beta \partial u \partial t + \gamma u \\ u(t, 0) & = f(t) \\ \partial u \partial x(t, 0) & = g(t) \\ u(0, x) & = h(x) \\ \partial u \partial t(0, x) & = v(x) \end{cases} \quad (1)$$

Taking  $Lu = \partial^2 u \partial x^2$  and  $Ru = \alpha \partial^2 u \partial t^2 + \beta \partial u \partial t + \gamma u$

We have :

$$Lu = Ru \quad (2)$$

Where  $L$  is an invertible operator in the Adomian sense and  $R$  the linear remainder.

Applying the laplace transform to the equation (2), we obtain :

$${}_x(Lu) = {}_x(Ru) \Leftrightarrow p_x^2(u) - pu(t, 0) - \partial u \partial x(t, 0) = {}_x(Ru) \quad (3)$$

$$p_x^2(u) = pf(t) + g(t) + {}_x(Ru) \quad (4)$$

Using the decomposition series for the linear term  $u(t, x)$  gives

$$p_{n \geq 0}^2(u_n) = pf(t) + g(t) + {}_{n \geq 0} {}_x(Ru_n) \quad (5)$$

This yields the following Adomian algorithm:

$$\begin{cases} p_x^2(u_0) = pf(t) + g(t) \\ p_x^2(u_{n+1}) = {}_x(Ru_n); n \geq 0 \end{cases} \quad (6)$$

Applying the laplace transform to the equation (2), we obtain :

$$\begin{cases} u_0(t, x) = {}_x^{-1} [1p^2(pf(t) + g(t))] \\ \dots \\ u_{n+1}(t, x) = {}_x^{-1} [1p_x^2(Ru_n)]; n \geq 0 \end{cases} \quad (7)$$

### 3 Algorithm of Laplace - ADM Convergence's

Considering the equation (1)

$$\begin{cases} \partial^2 u \partial x^2 & = \alpha \partial^2 u \partial t^2 + \beta \partial u \partial t + \gamma u \\ u(t, 0) & = f(t) \\ \partial u \partial x(t, 0) & = g(t) \\ u(0, x) & = h(x) \\ \partial u \partial t(0, x) & = v(x) \end{cases}$$

With  $(t, x) \in \Omega$  where  $\Omega = [0; +\infty[ \times [a, b]$

The application of the Laplace-ADM method gives

$$\begin{cases} u_0(t, x) = {}_x^{-1} [1p^2 (pf(t) + g(t))] \\ u_{n+1}(t, x) = {}_x^{-1} [1p_x^2 (Ru_n)]; n \geq 0 \end{cases}$$

**Let us suppose :**

·  $(H_1)$

$f$  is continuous then there is a real  $M$  so that

$$|f(t)| \leq M \text{ for all } t \in [0, T]$$

·  $(H_2)$

$g$  is continuous then there is a real  $M$  so that

$$|g(t)| \leq M' \text{ for all } t \in [0, T]$$

**However**

$R$  the linear remainder is continuous then there is a real  $\lambda > 0$  so that

$$\|Ru\| \leq \lambda \|u\|$$

**Indeed, we have :**

$$\begin{cases} |u_0| = |{}_x^{-1} [f(t)p] + {}_x^{-1} [g(t)p^2]| \\ \dots \dots \dots \\ |u_n| = |{}_x^{-1} [t(Ru_{n-1})p]|; n \geq 1 \end{cases}$$

There is a real  $x_0 \in \mathbb{R}_*^+$  so that  $\Re e(p) > x_0$ , we deduce the following system :

$$\begin{cases} |u_0| \leq |{}_x^{-1} [f(t)p]| + |{}_x^{-1} [g(t)p^2]| \\ \dots \dots \dots \\ |u_n| \leq |{}_x^{-1} [|Ru_{n-1}|p^2]|; n \geq 1 \end{cases}$$

$$\Rightarrow \begin{cases} |u_0| \leq M + M'b \\ \dots \dots \dots \\ |u_n| \leq |{}_x^{-1} [x(|Ru_{n-1}|)x_0^2]|; n \geq 1 \end{cases}$$

$$\Rightarrow \begin{cases} |u_0| \leq M + M'b \\ \dots \dots \dots \\ |u_n| \leq 1x_{0x}^{2-1} [x(|Ru_{n-1}^1|)]; n \geq 1 \end{cases}$$

$$\Rightarrow \begin{cases} |u_0| \leq M + M'b \\ \dots \dots \dots \\ |u_n| \leq \lambda x_0^2 ||u_{n-1}^1||; n \geq 1 \end{cases}$$

Step by step, we deduce :

$$\Rightarrow \begin{cases} |u_0| \leq M + M'b \\ \dots \quad \dots \quad \dots \\ |u_n| \leq (\lambda x_0^2)^n (M + M'b) ; n \geq 1 \end{cases}$$

With  $\lambda x_0^2 < 1$  and  $x_0 \neq \sqrt{\lambda}$ , we obtain

$$\Rightarrow \begin{cases} |u_0| \leq M + M'b \\ \dots \quad \dots \quad \dots \\ \sum_{n \geq 0} |u_n| \leq (M + M'b) x_0^2 x_0^2 - \lambda \end{cases}$$

Then the series  $\sum_{n \geq 0} u_n$  is convergent, therefore this algorithm is convergent.

## 4 Applications

### 4.1 Example : wave equation of microstrip antenna

In this example, we study the standard linear telegraph equation presented in wave equation of microstrip antenna equation is given as

$$\begin{cases} \partial^2 u \partial x^2 + \partial^2 u \partial t^2 + 2\partial u \partial t - u = 0 \\ \partial u \partial x (t, 0) = e^{-2t} \\ u(t, 0) = e^{-2t} \end{cases}$$

Taking  $Lu = \partial^2 u \partial x^2$ ,  $Ru = -\partial^2 u \partial t^2 - 2\partial u \partial t + u$ .

Where  $L$  is an invertible operator in the Adomian sense and  $R$  the linear remainder.

Applying the laplace transform to the equation (2), we obtain :

$${}_x(Lu) = {}_x(Ru) \quad (8)$$

$\Leftrightarrow$

$$p_x^2(u) - pu(t, 0) - \partial u \partial x(t, 0) = {}_x(-\partial^2 u \partial t^2 - 2\partial u \partial t + u) \quad (9)$$

$$(p^2 - 1)_x(u) = pe^{-2t} + e^{-2t} + {}_x(-\partial^2 u \partial t^2 - 2\partial u \partial t) \quad (10)$$

Using the decomposition series for the linear term  $u(t, x)$  gives

$$(p^2 - 1)_{n \geq 0} {}_x(u_n) = pe^{-2t} + e^{-2t} + {}_{n \geq 0} {}_x(-\partial^2 u_n \partial t^2 - 2\partial u_n \partial t) \quad (11)$$

We deduce the following Laplace-Adomian algorithm

$$\begin{cases} (p^2 - 1)_x(u_0) = pe^{-2t} + e^{-2t} \\ (p^2 - 1)_x(u_{n+1}) = {}_x(Ru_n) ; n \geq 0 \end{cases} \quad (12)$$

We obtain

$$\begin{cases} u_0(t, x) = \mathcal{L}_x^{-1} [1(p^2 - 1)(pe^{-2t} + e^{-2t})] \\ u_{n+1}(t, x) = \mathcal{L}_x^{-1} [1(p^2 - 1)_x (Ru_n)]; n \geq 0 \end{cases}$$

**Determinate**  $u_n(t, x)$ , for  $n \geq 0$

$$u_0(t, x) = \mathcal{L}_x^{-1} [1(p^2 - 1)(p + 1)e^{-2t}]$$

$$\Rightarrow u_0(t, x) = \mathcal{L}_x^{-1} [1(p - 1)e^{-2t}]$$

$$\Rightarrow u_0(t, x) = e^{x-2t}$$

$$u_1(t, x) = \mathcal{L}_t^{-1} [1(p^2 - 1)\mathcal{L}_t(R(u_0))]$$

$$\Rightarrow u_1(t, x) = \mathcal{L}_t^{-1} [1(p^2 - 1)(-4e^{x-2t} + 4e^{x-2t})]$$

$$\Rightarrow u_1(t, x) = 0$$

$$u_2(t, x) = \mathcal{L}_t^{-1} [1(p^2 - 1)\mathcal{L}_t(R(u_1))]$$

$$\Rightarrow u_2(t, x) = 0$$

In recursive way, we deduce

$$u_n(t, x) = 0 \text{ for all } n \geq 1$$

Then

$$u(t, x) =_{n \geq 0} u_n(t, x) = e^{x-2t}$$

The exact solution of model is

$$u(t, x) = e^{x-2t}$$

## 4.2 Example : wave equation of microstrip antenna

In this example, we study the standard linear telegraph equation is given as

$$\begin{cases} \partial^2 u \partial x^2 & = \partial^2 u \partial t^2 + 4\partial u \partial t + 4u \\ \partial u \partial t(0, x) & = -2 \\ u(0, x) & = 1 + e^{2x} \end{cases}$$

Taking  $Lu = \partial^2 u \partial t^2$ ,  $Ru = \partial^2 u \partial x^2 - 4\partial u \partial t - 4u$ .

Where  $L$  is an invertible operator in the Adomian sense and  $R$  the linear remainder.

Applying the laplace transform to the equation (2), we obtain :

$${}_t(Lu) = {}_t(Ru) \Leftrightarrow p_t^2(u) - pu(0, x) - \partial u \partial x(0, x) = {}_t(\partial^2 u \partial x^2 - 4\partial u \partial t - 4u) \quad (13)$$

$$p_x^2(u) = p(1 + e^{2x}) - 2 + {}_x(\partial^2 u \partial x^2 - 4\partial u \partial t - 4u) \quad (14)$$

Using the decomposition series for the linear term  $u(t, x)$  gives

$$p_{n \geq 0}^2(u_n) = p - 2 + pe^{2x} + {}_{n \geq 0}{}_x(\partial^2 u_n \partial x^2 - 4\partial u_n \partial t - 4u) \quad (15)$$

$$p_{n \geq 0}^2(u_n) = p - 2 + pe^{2x} + {}_{n \geq 0}{}_x(\partial^2 u_n \partial x^2 - 4\partial u_n \partial t - 4u) \quad (16)$$

We deduce the following Laplace-Adomian algorithm

$$\begin{cases} p_x^2(u_0) = p - 2 + pe^{2x} \\ p_x^2(u_{n+1}) = {}_x(Ru_n); n \geq 0 \end{cases} \quad (17)$$

We obtain

$$\begin{cases} u_0(t, x) = {}_x^{-1} [1p^2(p - 2 + pe^{2x})] \\ u_{n+1}(t, x) = {}_x^{-1} [1p_x^2(Ru_n)]; n \geq 0 \end{cases}$$

**Determinate**  $u_n(t, x)$ , for  $n \geq 0$

$$\begin{aligned} u_0(t, x) &= {}_t^{-1} [1p^2(p - 2 + pe^{2x})] \\ \Rightarrow u_0(t, x) &= {}_t^{-1} [1p(1 + e^{2x}) - 2p^2] \\ \Rightarrow u_0(t, x) &= e^{2x} - 2t + 1 \end{aligned}$$

$$\begin{aligned} u_1(t, x) &= {}_t^{-1} [1p_t^2(R(u_0))] \\ \Rightarrow u_1(t, x) &= {}_t^{-1} [1p^2 [{}_t(-8t + 4)]] \end{aligned}$$

$$\Rightarrow u_1(t, x) = {}_t^{-1} (-8p^4 + 4p^3) = -8t^3 3! + 4t 2!$$

$$\Rightarrow u_1(t, x) = -(2t)^3 3! + (2t)^2 2!$$

$$\begin{aligned} u_2(t, x) &= {}_t^{-1} [1p_t^2(R(u_1))] \\ \Rightarrow u_2(t, x) &= {}_t^{-1} [1p^2 [{}_t(-16t - 24t^2 - 32t^3 3!)]] \end{aligned}$$

$$\Rightarrow u_2(t, x) = {}_t^{-1} (-161p^4 - 481p^5 - 321p^6)$$

$$\Rightarrow u_2(t, x) = -2(2t)^3 3! - 3(2t)^4 4! - (2t)^5 5!$$

$$u_3(t, x) = {}_t^{-1} [1p^2_t (R(u_2))] \\ \Rightarrow u_3(t, x) = {}_t^{-1} [1p^2 [ {}_t (2^6 t^2 2! + 2^8 t^3 3! + 10 \times 2^5 t^4 4! + 2^7 t^5 5!)]]$$

$$\Rightarrow u_3(t, x) = {}_t^{-1} (2^6 1p^5 + 2^8 1p^6 + 5 \times 2^6 1p^7 + 2^7 1p^8)$$

$$\Rightarrow u_3(t, x) = 4(2t)^4 4! + 8(2t)^5 5! + 5(2t)^6 6! + (2t)^7 7!$$

$$u_4(t, x) = {}_t^{-1} [1p^2_t (R(u_3))] \\ \Rightarrow u_4(t, x) = {}_t^{-1} [1p^2 [ {}_t (-\frac{32}{315} t^7 - \frac{112}{45} t^6 - \frac{96}{5} t^5 - \frac{160}{3} t^4 - \frac{128}{3} t^3) ]] \\ \Rightarrow u_4(t, x) = {}_t^{-1} [1p^2 [ {}_t (-2^8 t^3 3! - 5 \times 2^8 t^4 4! - 9 \times 2^8 t^5 5! - 7 \times 2^8 t^6 6! - 2^9 t^7 7!)]] \\ \Rightarrow u_4(t, x) = {}_t^{-1} (-2^8 1p^6 - 5 \times 2^8 1p^7 - 9 \times 2^8 1p^8 - 7 \times 2^8 1p^9 - 2^9 1p^{10})$$

$$\Rightarrow u_4(t, x) = -8(2t)^5 5! - 20(2t)^6 6! - 18(2t)^7 7! - 7(2t)^8 8! - (2t)^9 9!$$

Step by step, we deduce

$${}_{n \geq 0} u_n(t, x) = e^{2x} + {}_{n \geq 0} (-2t)^k n!$$

Then

$$u(t, x) = {}_{n \geq 0} u_n(t, x) = e^{2x} + e^{-2t}$$

The exact solution of model is

$$u(t, x) = e^{2x} + e^{-2t}$$

## 5 Conclusion

Laplace's Adomian numerical method allowed us to solve some linear partial differential equations by modelling the standard telegraph equation. It is therefore a very powerful numerical analysis tool to solve this type of problem, this method accelerates convergence to the solution. Our study was limited to the linear models of telegraph non-homogeneous reaction, a study of these models in non-homogeneous cases would be an important contribution to the understanding of these models.

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