

Resolution of the standard telegraph equation by the Laplace-Adomian method

Abstract

In this paper, we research the solution of the standard telegraph equation by the Laplace-Adomian method. The Laplace-Adomian method is based on the combination of Laplace transform and the Adomian decomposition method.

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Key word: Telegraph equation, Laplace transform, ADM method.

1 Introduction

In this article, we study the general solution of the standard telegraph equations by the method of Laplace-Adomian. The standard telegraph equation is an important equation arises in the propagation of electrical signals along a telegraph line, taking into consideration the inductance, capacitance and conductance of the cable. However the method of Laplace-Adomian is a numerical method based on the combination of the Laplace Transform and Adomian decomposition method.

2 The numerical Laplace-Adomian method

The standard telegraph equation is a partial differential equation given by :

$$\partial u^2 \partial x^2 = \alpha \partial^2 u \partial t^2 + \beta \partial u \partial t + \gamma u$$

where $u = u(t, x)$ is the resistance, and a, b and c are constants related to the inductance, capacitance and conductance of the cable respectively.

Let us consider the following functional equation:

$$\begin{cases} \partial u^2 \partial x^2 & = \alpha \partial^2 u \partial t^2 + \beta \partial u \partial t + \gamma u \\ u(t, 0) & = f(t) \\ \partial u \partial x(t, 0) & = g(t) \\ u(0, x) & = h(x) \\ \partial u \partial t(0, x) & = v(x) \end{cases} \quad (1)$$

Taking $Lu = \partial u^2 \partial x^2$ and $Ru = \alpha \partial^2 u \partial t^2 + \beta \partial u \partial t + \gamma u$

We have :

$$Lu = Ru \quad (2)$$

Where L is an invertible operator in the Adomian sense and R the linear remainder.

Applying the laplace transform to the equation (2), we obtain :

$${}_x(Lu) = {}_x(Ru) \Leftrightarrow p_x^2(u) - pu(t, 0) - \partial u \partial x(t, 0) = {}_x(Ru) \quad (3)$$

$$p_x^2(u) = pf(t) + g(t) + {}_x(Ru) \quad (4)$$

Using the decomposition series for the linear term $u(t, x)$ gives

$$p_{n \geq 0}^2(u_n) = pf(t) + g(t) + {}_{n \geq 0}{}_x(Ru_n) \quad (5)$$

This yields the following Adomian algorithm:

$$\begin{cases} p_x^2(u_0) = pf(t) + g(t) \\ p_x^2(u_{n+1}) = {}_x(Ru_n); n \geq 0 \end{cases} \quad (6)$$

Applying the laplace transform to the equation (2), we obtain :

$$\begin{cases} u_0(t, x) = {}_x^{-1} [1p^2(pf(t) + g(t))] \\ u_{n+1}(t, x) = {}_x^{-1} [1p_x^2(Ru_n)]; n \geq 0 \end{cases} \quad (7)$$

3 Algorithm of Laplace - ADM Convergence's

Considering the equation (1)

$$\begin{cases} \partial u^2 \partial x^2 & = \alpha \partial^2 u \partial t^2 + \beta \partial u \partial t + \gamma u \\ u(t, 0) & = f(t) \\ \partial u \partial x(t, 0) & = g(t) \\ u(0, x) & = h(x) \\ \partial u \partial t(0, x) & = v(x) \end{cases}$$

With $(t, x) \in \Omega$ where $\Omega = [0; +\infty[\times [a, b]$

The application of the Laplace-ADM method gives

$$\begin{cases} u_0(t, x) = {}_x^{-1} [1p^2 (pf(t) + g(t))] \\ u_{n+1}(t, x) = {}_x^{-1} [1p^2_x (Ru_n)]; n \geq 0 \end{cases}$$

Let us suppose :

· (H_1)

f is continuous then there is a real M so that

$$|f(t)| \leq M \text{ for all } t \in [0, T]$$

· (H_2)

g is continuous then there is a real M so that

$$|g(t)| \leq M' \text{ for all } t \in [0, T]$$

However

R the linear remainder is continuous then there is a real $\alpha > 0$ so that

$$\|Ru\| \leq \alpha \|u\|$$

Indeed, we have :

$$\begin{cases} |u_0| = |{}_x^{-1} [f(t)p] + {}_x^{-1} [g(t)p^2]| \\ |u_n| = |{}_x^{-1} [t(Ru_{n-1})p]|; n \geq 1 \end{cases}$$

There is a real $x_0 \in \mathbb{R}_*^+$ so that $\Re e(p) > x_0$, we deduce the following system :

$$\begin{aligned} & \begin{cases} |u_0| \leq |{}_x^{-1} [f(t)p]| + |{}_x^{-1} [g(t)p^2]| \\ |u_n| \leq |{}_x^{-1} [t(Ru_{n-1})p]|; n \geq 1 \end{cases} \\ \Rightarrow & \begin{cases} |u_0| \leq M + M'b \\ |u_n| \leq |{}_x^{-1} [x(|Ru_{n-1}|)x_0^2]|; n \geq 1 \end{cases} \\ \Rightarrow & \begin{cases} |u_0| \leq M + M'b \\ |u_n| \leq 1x_{0x}^{2-1} [x(|Ru_{n-1}^1|)]; n \geq 1 \end{cases} \\ \Rightarrow & \begin{cases} |u_0| \leq M + M'b \\ |u_n| \leq \alpha x_0^2 ||u_{n-1}^1||; n \geq 1 \end{cases} \end{aligned}$$

Step by step, we deduce :

$$\begin{aligned} \Rightarrow & \begin{cases} |u_0| \leq M + M'b \\ |u_n| \leq (\alpha x_0^2)^n (M + M'b); n \geq 1 \end{cases} \\ \Rightarrow & \begin{cases} |u_0| \leq M + M'b \\ \sum_{n \geq 0} |u_n| \leq (M + M'b) x_0^2 x_0^2 - \alpha \end{cases} \end{aligned}$$

Then the series $\sum_{n \geq 0} u_n$ is convergent, therefore this algorithm is convergent.

4 Applications

4.1 Example : wave equation of microstrip antenna

In this example, we study the standard linear telegraph equation presented in wave equation of microstrip antenna equation is given as

$$\begin{cases} \partial u^2 \partial x^2 + \partial^2 u \partial t^2 + 2 \partial u \partial t - u & = & 0 \\ \partial u \partial x (t, 0) & = & e^{-2t} \\ u(t, 0) & = & e^{-2t} \end{cases}$$

Taking $Lu = \partial u^2 \partial x^2$, $Ru = -\partial^2 u \partial t^2 - 2 \partial u \partial t + 2u$.

Where L is an invertible operator in the Adomian sense and R the linear remainder.

Applying the laplace transform to the equation (2), we obtain :

$${}_x(Lu) = {}_x(Ru) \tag{8}$$

\Leftrightarrow

$$p_x^2(u) - pu(t, 0) - \partial u \partial x(t, 0) = {}_x(-\partial^2 u \partial t^2 - 2 \partial u \partial t + u) \tag{9}$$

$$(p^2 - 1)_x(u) = pe^{-2t} + e^{-2t} + {}_x(-\partial^2 u \partial t^2 - 2 \partial u \partial t) \tag{10}$$

Using the decomposition series for the linear term $u(t, x)$ gives

$$(p^2 - 1)_{n \geq 0} {}_x(u_n) = pe^{-2t} + e^{-2t} + {}_{n \geq 0} {}_x(-\partial^2 u_n \partial t^2 - 2 \partial u_n \partial t) \tag{11}$$

We deduce the following Laplace-Adomian algorithm

$$\begin{cases} (p^2 - 1)_x(u_0) = pe^{-2t} + e^{-2t} \\ (p^2 - 1)_x(u_{n+1}) = {}_x(Ru_n); n \geq 0 \end{cases} \tag{12}$$

We obtain

$$\begin{cases} u_0(t, x) = {}_x^{-1} [1(p^2 - 1)(pe^{-2t} + e^{-2t})] \\ u_{n+1}(t, x) = {}_x^{-1} [1(p^2 - 1)_x(Ru_n)]; n \geq 0 \end{cases}$$

Determinate $u_n(t, x)$, for $n \geq 0$

$$u_0(t, x) = {}_x^{-1} [1(p^2 - 1)(p + 1)e^{-2t}]$$

$$\Rightarrow u_0(t, x) = {}_x^{-1} [1(p - 1)e^{-2t}]$$

$$\Rightarrow u_0(t, x) = e^{x-2t}$$

$$u_1(t, x) = \mathcal{L}_t^{-1} [1(p^2 - 1)\mathcal{L}_t(R(u_0))]$$

$$\Rightarrow u_1(t, x) = \mathcal{L}_t^{-1} [1(p^2 - 1)(-4e^{x-2t} + 4e^{x-2t})]$$

$$\Rightarrow u_1(t, x) = 0$$

$$u_2(t, x) = \mathcal{L}_t^{-1} [1(p^2 - 1)\mathcal{L}_t(R(u_1))] \\ \Rightarrow u_2(t, x) = 0$$

In recursive way, we deduce

$$u_n(t, x) = 0 \text{ for all } n \geq 1$$

Then

$$u(t, x) =_{n \geq 0} u_n(t, x) = e^{x-2t}$$

The exact solution of model is

$$u(t, x) = e^{x-2t}$$

4.2 Example : wave equation of microstrip antenna

In this example, we study the standard linear telegraph equation is given as

$$\begin{cases} \partial u^2 \partial x^2 & = \partial^2 u \partial t^2 + 4 \partial u \partial t + 4u \\ \partial u \partial t(0, x) & = -2 \\ u(0, x) & = 1 + e^{2x} \end{cases}$$

Taking $Lu = \partial u^2 \partial t^2$, $Ru = \partial^2 u \partial x^2 - 4 \partial u \partial t - 4u$.

Where L is an invertible operator in the Adomian sense and R the linear remainder.

Applying the laplace transform to the equation (2), we obtain :

$${}_t(Lu) = {}_t(Ru) \Leftrightarrow p_t^2(u) - pu(0, x) - \partial u \partial x(0, x) = {}_t(\partial^2 u \partial x^2 - 4 \partial u \partial t - 4u) \quad (13)$$

$$p_x^2(u) = p(1 + e^{2x}) - 2 + {}_x(\partial^2 u \partial x^2 - 4 \partial u \partial t - 4u) \quad (14)$$

Using the decomposition series for the linear term $u(t, x)$ gives

$$p_{n \geq 0}^2(u_n) = p - 2 + pe^{2x} + {}_{n \geq 0} {}_x(\partial^2 u_n \partial x^2 - 4 \partial u_n \partial t - 4u) \quad (15)$$

We deduce the following Laplace-Adomian algorithm

$$\begin{cases} p_x^2(u_0) = p - 2 + pe^{2x} \\ p_x^2(u_{n+1}) =_x (Ru_n); n \geq 0 \end{cases} \quad (16)$$

We obtain

$$\begin{cases} u_0(t, x) =_x^{-1} [1p^2(p - 2 + pe^{2x})] \\ u_{n+1}(t, x) =_x^{-1} [1p_x^2(Ru_n)]; n \geq 0 \end{cases}$$

Determinate $u_n(t, x)$, for $n \geq 0$

$$\begin{aligned} u_0(t, x) &= _t^{-1} [1p^2(p - 2 + pe^{2x})] \\ \Rightarrow u_0(t, x) &= _t^{-1} [1p(1 + e^{2x}) - 2p^2] \\ \Rightarrow u_0(t, x) &= e^{2x} - 2t + 1 \end{aligned}$$

$$\begin{aligned} u_1(t, x) &= _t^{-1} [1p_t^2(R(u_0))] \\ \Rightarrow u_1(t, x) &= _t^{-1} [1p^2[_t(8t + 4)]] \\ \Rightarrow u_1(t, x) &= _t^{-1} (8p^4 + 4p^3) = 8t^33! + 4t2! \\ \Rightarrow u_1(t, x) &= (2t)^33! + (2t)^22! \end{aligned}$$

$$\begin{aligned} u_2(t, x) &= _t^{-1} [1p_t^2(R(u_1))] \\ \Rightarrow u_2(t, x) &= _t^{-1} [1p^2[_t(-16t - 24t^2 - 32t^33!)]] \\ \Rightarrow u_2(t, x) &= _t^{-1} (-161p^4 - 481p^5 - 321p^6) \\ \Rightarrow u_2(t, x) &= -2(2t)^33! - 3(2t)^44! - (2t)^55! \end{aligned}$$

$$\begin{aligned} u_3(t, x) &= _t^{-1} [1p_t^2(R(u_2))] \\ \Rightarrow u_3(t, x) &= _t^{-1} [1p^2[_t(2^6t^22! + 2^8t^33! + 10 \times 2^5t^44! + 2^7t^55!)]] \\ \Rightarrow u_3(t, x) &= _t^{-1} (2^61p^5 + 2^81p^6 + 5 \times 2^61p^7 + 2^71p^8) \\ \Rightarrow u_3(t, x) &= 4(2t)^44! + 8(2t)^55! + 5(2t)^66! + (2t)^77! \end{aligned}$$

$$\begin{aligned} u_4(t, x) &= _t^{-1} [1p_t^2(R(u_3))] \\ \Rightarrow u_4(t, x) &= _t^{-1} [1p^2[_t(-\frac{32}{315}t^7 - \frac{112}{45}t^6 - \frac{96}{5}t^5 - \frac{160}{3}t^4 - \frac{128}{3}t^3)]] \\ \Rightarrow u_4(t, x) &= _t^{-1} [1p^2[_t(-2^8t^33! - 5 \times 2^8t^44! - 9 \times 2^8t^55! - 7 \times 2^8t^66! - 2^9t^77!)]] \\ \Rightarrow u_4(t, x) &= _t^{-1} (-2^81p^6 - 5 \times 2^81p^7 - 9 \times 2^81p^8 - 7 \times 2^81p^9 - 2^91p^{10}) \\ \Rightarrow u_4(t, x) &= -8(2t)^55! - 20(2t)^66! - 18(2t)^77! - 7(2t)^88! - (2t)^99! \end{aligned}$$

Step by step, we deduce

$${}_{n \geq 0} u_n(t, x) = e^{2x} + {}_{n \geq 0} (-2t)^k n!$$

Then

$$u(t, x) = {}_{n \geq 0} u_n(t, x) = e^{2x} + e^{-2t}$$

The exact solution of model is

$$u(t, x) = e^{2x} + e^{-2t}$$

5 Conclusion

Laplace's Adomian numerical method allowed us to solve some linear partial differential equations by modelling the standard telegraph equation. It is therefore a very powerful numerical analysis tool to solve this type of problem, this method accelerates convergence to the solution. Our study was limited to the linear models of telegraph non-homogeneous reaction, a study of these models in non-homogeneous cases would be an important contribution to the understanding of these models.

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